

MDOF SYSTEMS WITH DAMPING

MDOF Systems with hysteretic damping- general case

Free vibration solution:

$$[M]\{\ddot{x}\} + ([K] + i[D])\{x\} = \{0\}$$

Assume a solution in the form of:

$$\{x\} = \{X\}e^{i\lambda t}$$

Here λ can be a complex number. The solution here is like the undamped case. However, both eigenvalues and Eigenvector matrices are complex.

The eigensolution has the orthogonal properties as:

$$[\Psi]^T [M] [\Psi] = [m_r]; \quad [\Psi]^T [K + iD] [\Psi] = [k_r]$$

The modal mass and stiffness parameters are complex.

MDOF Systems with hysteretic damping- general case

Again, the following relation is valid:

$$\lambda_r^2 = \frac{k_r}{m_r} = \omega_r^2 (1 + i\eta_r)$$

A set of mass-normalized eigenvectors can be defined as:

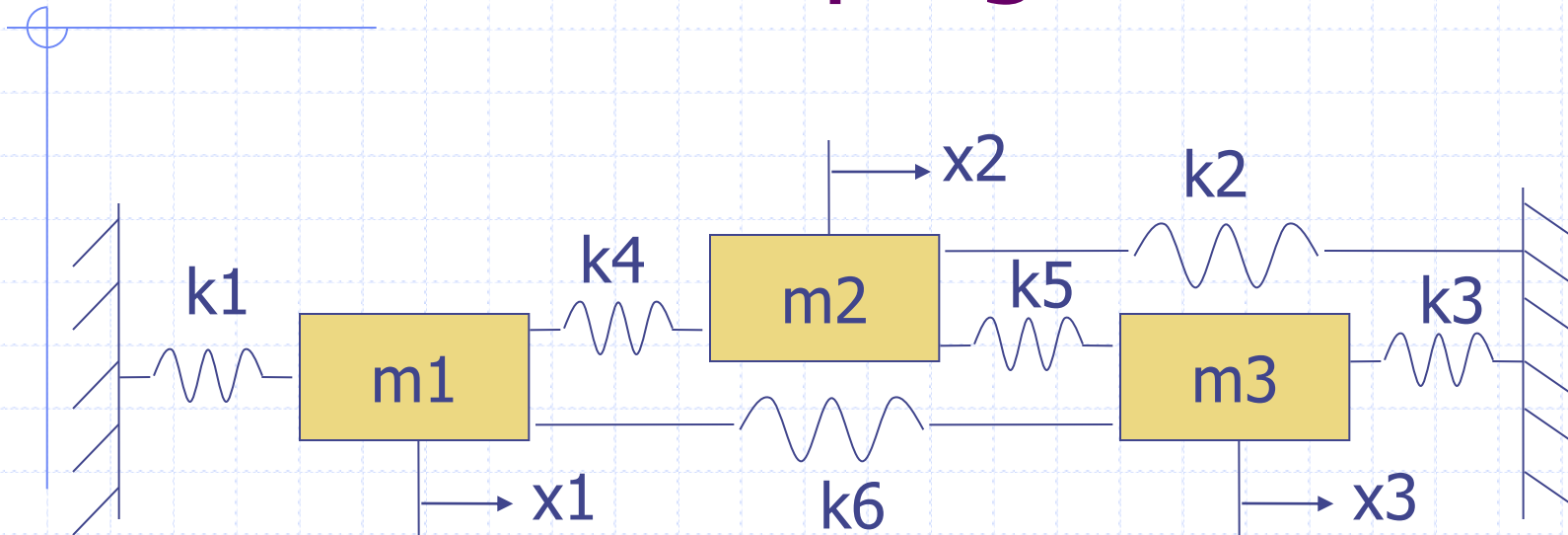
$$\{\phi\}_r = (m_r)^{-1/2} \{\psi\}_r$$

What is the interpretation of complex mode shapes?

The phase angle in undamped is either 0 or 180.

Here the phase angle may take any value.

Numerical Example with structural damping



$$m_1 = 0.5 \text{ Kg}$$

$$m_2 = 1.0 \text{ Kg}$$

$$m_3 = 1.5 \text{ Kg}$$

$$k_1 = k_2 = k_3 = k_4 = k_5 = k_6 = 1000 \text{ N/m}$$

Undamped

$$[M] = \begin{bmatrix} .5 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1.5 \end{bmatrix} \quad [K] = \begin{bmatrix} 3 & -1 & -1 \\ -1 & 3 & -1 \\ -1 & -1 & 3 \end{bmatrix}$$

Using command `[V,D]=eig(k,M)` in MATLAB

$$[\omega_r^2] = \begin{bmatrix} 950 & 0 & 0 \\ 0 & 3352 & 0 \\ 0 & 0 & 6698 \end{bmatrix} \quad [\Phi] = \begin{bmatrix} .464 & -.218 & -1.318 \\ .536 & -.782 & .318 \\ .635 & .493 & .142 \end{bmatrix}$$

Proportional Structural Damping

Assume proportional structural damping as:

$$d_j = 0.05k_j, \quad j = 1, \dots, 6$$

$$[\lambda_r^2] = \begin{bmatrix} 950(1 + .05i) & 0 & 0 \\ 0 & 3352(1 + .05i) & 0 \\ 0 & 0 & 6698(1 + .05i) \end{bmatrix}$$

$$[\Phi] = \begin{bmatrix} .464(0^\circ) & .218(180^\circ) & 1.318(180^\circ) \\ .536(0^\circ) & .782(180^\circ) & .318(0^\circ) \\ .635(0^\circ) & .493(0^\circ) & .142(0^\circ) \end{bmatrix}$$

Non-Proportional Structural Damping

Assume non-proportional structural damping as:

$$d_1 = 0.3k_1$$

$$d_j = 0, \quad j = 2, \dots, 6$$

$$[\lambda_r^2] = \begin{bmatrix} 957(1 + .067i) & 0 & 0 \\ 0 & 3354(1 + .042i) & 0 \\ 0 & 0 & 6690(1 + .078i) \end{bmatrix}$$

$$[\Phi] = \begin{bmatrix} .463(-5.5^\circ) & .217(173^\circ) & 1.321(181^\circ) \\ .537(0^\circ) & .784(181^\circ) & .316(-6.7^\circ) \\ .636(0^\circ) & .492(-1.3^\circ) & .142(-3.1^\circ) \end{bmatrix}$$

Non-Proportional Structural Damping

- ◆ Each mode has a different damping factor.
- ◆ All eigenvectors arguments for undamped and proportional damp cases are either 0 or 180.
- ◆ All eigenvectors arguments for non-proportional case are within 10 degree of 0 or 180 (the modes are almost real).

Exercise: Repeat the problem with
 $m_1=1\text{Kg}$, $m_2=0.95\text{ Kg}$, $m_3=1.05\text{ Kg}$
 $k_1=k_2=k_3=k_4=k_5=k_6=1000\text{ N/m}$

FRF Characteristics (Hysteretic Damping)

Again, one can write:

$$([K] + i[D] - \omega^2[M])\{X\}e^{i\omega t} = \{F\}e^{i\omega t}$$

The receptance matrix can be found as:

$$H(\omega) = ([K] + i[D] - \omega^2[M])^{-1} = [\Phi][\lambda_r^2 - \omega^2][\Phi]^T$$

FRF elements can be extracted:

$$H_{jk}(\omega) = \sum_{r=1}^N \frac{\phi_{jr}\phi_{kr}}{\omega_r^2 - \omega^2 + i\eta_r\omega_r^2}$$

or

$$H_{jk}(\omega) = \sum_{r=1}^N \frac{{}_r A_{jk}}{\omega_r^2 - \omega^2 + i\eta_r\omega_r^2}$$

Modal Constant

MDOF Systems with viscous damping- general case

The general equation of motion for this case can be written as:

$$[M]\{\ddot{x}\} + [C]\{\dot{x}\} + [K]\{x\} = \{f\}$$

Consider the zero excitation to determine the natural frequencies and mode shapes of the system:

$$\{x\} = \{X\}e^{st}$$

This leads to:

$$([M]s^2 + [C]s + [K])\{X\} = \{0\}$$

This is a complex eigenproblem. In this case, there are $2N$ eigenvalues but they are in complex conjugate pairs.

MDOF Systems with viscous damping- general case

$$\begin{cases} s_r, s_r^* \\ \{\psi\}_r, \{\psi\}_r^* \end{cases} \quad r = 1, \dots, N$$

It is customary to express each eigenvalues as:

$$s_r = \omega_r (-\zeta_r + i\sqrt{1 - \zeta_r^2})$$

Next, consider the following equation:

$$(s_r^2 [M] + s_r [C] + [K]) \{\psi\}_r = \{0\}$$

Then, pre-multiply by $\{\psi\}_q^H$:

$$\{\psi\}_q^H (s_r^2 [M] + s_r [C] + [K]) \{\psi\}_r = \{0\}$$

*

MDOF Systems with viscous damping- general case

A similar expression can be written for $\{\psi\}_q$:

$$(s_q^2[M] + s_q[C] + [K])\{\psi\}_q = \{0\}$$

This can be transposed-conjugated and then multiply by $\{\psi\}_r$

$$\{\psi\}_q^H (s_q^2[M] + s_q[C] + [K])\{\psi\}_r = \{0\} \quad **$$

Subtract equation * from **, to get:

$$(s_r^2 - s_q^2)\{\psi\}_q^H [M]\{\psi\}_r + (s_r - s_q)\{\psi\}_q^H [C]\{\psi\}_r = \{0\}$$

This leads to the first orthogonality equations:

$$(s_r + s_q)\{\psi\}_q^H [M]\{\psi\}_r + \{\psi\}_q^H [C]\{\psi\}_r = \{0\} \quad (1)$$

MDOF Systems with viscous damping- general case

Next, multiply equation (*) by s_q and (**) by s_r :

$$s_r s_q \{\psi\}_q^H [M] \{\psi\}_r - \{\psi\}_q^H [K] \{\psi\}_r = \{0\} \quad (2)$$

Equations (1) and (2) are the orthogonality conditions:
If we use the fact that the modes are pair, then

$$s_q = \omega_r (-\zeta_r - i\sqrt{1-\zeta_r^2})$$

$$\{\psi\}_q = \{\psi\}_r^*$$

MDOF Systems with viscous damping- general case

Inserting these two into equations (1) and (2):

$$2\omega_r \zeta_r = \frac{\{\psi\}_r^H [C] \{\psi\}_r}{\{\psi\}_r^H [M] \{\psi\}_r} = \frac{c_r}{m_r}$$

$$\omega_r^2 = \frac{\{\psi\}_r^H [K] \{\psi\}_r}{\{\psi\}_r^H [M] \{\psi\}_r} = \frac{k_r}{m_r}$$

Where m_r , k_r , c_r are modal mass, stiffness and damping.

FRF Characteristics (Viscous Damping)

The response solution is:

$$\{X\} = ([K] + i\omega[C] - \omega^2[M])^{-1}\{F\}$$

We are seeking to a similar series expansion similar to the undamped case.

To do this, we define a new vector $\{u\}$:

$$\{u\} = \begin{Bmatrix} x \\ \dot{x} \end{Bmatrix}_{2N \times 1}$$

We write the equation of motion as:

$$[C : M]_{N \times 2N} \{\dot{u}\}_{2N \times 1} + [K : 0]\{u\} = \{0\}_{N \times 1}$$

FRF Characteristics (Viscous Damping)

This is N equations and 2N unknowns. We add an identity Equation as:

$$[M : 0]\{\dot{u}\} + [0 : -M]\{u\} = \{0\}$$

Now, we combine these two equations to get:

$$\begin{bmatrix} C & M \\ M & 0 \end{bmatrix}\{\dot{u}\} + \begin{bmatrix} K & 0 \\ 0 & -M \end{bmatrix}\{u\} = \{0\}$$

Which can be simplified to:

$$[A]\{\dot{u}\} + [B]\{u\} = \{0\}$$

FRF Characteristics (Viscous Damping)

Equation (3) is in a standard eigenvalue form. Assuming a trial solution in the form of $\{u\} = \{U\}e^{st}$

$$(s_r[A] + [B])\{\theta\}_r = \{0\} \quad r = 1, \dots, 2N$$

The orthogonality properties can be stated as:

$$[\theta]^T [A] [\theta] = [a_r]$$

$$[\theta]^T [B] [\theta] = [b_r]$$

With the usual characteristics:

$$s_r = \frac{b_r}{a_r} \quad r = 1, \dots, 2N$$

FRF Characteristics (Viscous Damping)

Let's express the forcing vector as:

$$\{P\}_{2N \times 1} = \begin{Bmatrix} F \\ 0 \end{Bmatrix}$$

Now using the previous series expansion:

$$\begin{Bmatrix} X \\ i\omega X \end{Bmatrix}_{2N \times 1} = \sum_{r=1}^{2N} \frac{\{\theta\}_r^T \{P\} \{\theta\}_r}{a_r (i\omega - s_r)}$$

And because the eigenvalues and vectors occur in complex conjugate pair:

$$\begin{Bmatrix} X \\ i\omega X \end{Bmatrix}_{2N \times 1} = \sum_{r=1}^N \frac{\{\theta\}_r^T \{P\} \{\theta\}_r}{a_r (i\omega - s_r)} + \frac{\{\theta\}_r^H \{P\} \{\theta^*\}_r}{a_r^* (i\omega - s_r^*)}$$

FRF Characteristics (Viscous Damping)

Now the receptance frequency response function H_{jk}

Resulting from a single force F_k and response parameter X_j

$$H_{jk}(\omega) = \sum_{r=1}^N \frac{\theta_{jr} \theta_{kr}}{a_r (\omega_r \zeta_r + i(\omega - \omega_r \sqrt{1 - \zeta_r^2}))} + \frac{\theta_{jr}^* \theta_{kr}^*}{a_r^* (\omega_r \zeta_r + i(\omega + \omega_r \sqrt{1 - \zeta_r^2}))}$$

or

$$H_{jk}(\omega) = \sum_{r=1}^N \frac{{}_r R_{jk} + i(\omega / \omega_r)({}_r S_{jk})}{\omega_r^2 - \omega^2 + 2i\omega\omega_r \zeta_r}$$

Where:

$$\{{}_r R_k\} = 2(\zeta_r \operatorname{Re}\{{}_r G_k\} - \operatorname{Im}\{{}_r G_k\} \sqrt{1 - \zeta_r^2})$$

$$\{{}_r S_k\} = 2 \operatorname{Re}\{{}_r G_k\}$$

$$\{{}_r G_k\} = (\theta_{kr} / a_r) \{\theta\}_r$$